Structure and Raman spectrum of BiSr₂CaCu₂O₇ (Bi₁₂₁₂) superconductor

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Abstract

 $BiSr_2CaCu_2O_7$ (Bi_{1212}) is prepared and characterized by various methods in particular by Raman spectroscopy. The material is found to be very anisotrope. The identity between the polarized Raman spectra along XX and YY axes is in agreement with a tetragonal structure P4/mmm. The collapse of the stretching (Bi)O(Cu) from two Raman bands in Bi_{2212} into only one band in Bi_{1212} reflects a vertically aligned bond ordering, similar to that observed when converting from YBa₂Cu₄O₈ into YBa₂Cu₃O₇ and from $Tl_2Ba_2CaCu_2O_8$ into TlBa₂CaCu₂O₇.

1 - Introduction

The easy structural conversion by thermal treatment⁽¹⁻³⁾ from YBa₂Cu₄O₈ (Y₁₂₄) into $YBa_2Cu_3O_7$ (Y₁₂₃) and from $T\ell_2Ba_2CaCu_2O_8$ $(T\ell_{2212})^{(4)}$ into $T\ell Ba_2 Ca Cu_2 O_7$ (Tl_{1212}) and their analogy with the bismuth cuprate superconductors suggest that the same conversion could happen from Bi₂Sr₂CaCu₂O₈ (Bi₂₂₁₂) into $BiSr_2CaCu_2O_7$ (Bi_{1212}). The Bi series also have their interest in practical applications. The working temperature of a superconducting device must be at least 25% lower than its T_c , due to losses of various nature $^{(5)}$. So, yttrium series have too low T_c for liquid nitrogen temperature cooling devices. On the other hand, thallium is rather toxic; its handling is thus delicate, especially in the synthesis phase. The bismuth series do not present such disadvantages. Nevertheless if the two or three BiO layers compounds are well studied, it is not the case for one BiO layer series.

Recently $(Bi_{0.5}Cu_{0.5})Sr_2YCa_2O_7$ was obtained by Ehmann *et al.*⁽⁶⁾ and identified as having P4/mmm structure.

In this paper, we present our study on $BiSr_2CaCu_2O_7$ (Bi_{1212}) containing only one BiO chain/layer. The material was characterized by several physical methods especially by Raman spectroscopy, this latter technique has been

revealed to be very helpful in the structural study of superconductors.

2 - Experimental

Bi₁₂₁₂ was prepared by sintering a calcinated lyophilized precursor as follows. Fine drops of about one hundred um in diameter of a stoichiometric aqueous solution (15 g/ ℓ) of nitrates (Bi, Sr, Ca, Cu) were pulverized in liquid nitrogen. After detente, the frozen powder was then liophylized progressively from -30 to $+80^{\circ}$ C during 24 h under 10-3 atm. pressure. A further calcination at 520°C during 18 h is managed to decompose the nitrates and to avoid the rehydration. This temperature was respected to prevent the eutectic (0.9Bi₂O₃-0.1CuO) which melts at 570°C and induces irreversible microstructural defects^(7,8). After pelleting under uni-directional pressure, the sample was sintered at 800°C in air for 100 h followed by a final annealing at 500° C in an oxygen flow.

The Raman spectra were recorded on a microspectrometer Dilor/Omars equipped with a multichannel photodiodes detector and an ion/argon laser source Spectra-Physics, model 2016. The spatial resolution was around 1 μ m², the laser power was some microwatts and the spectral resolution was 3 cm⁻¹.



Fig.1 - Polarized Raman spectra of Bi₁₂₁₂.

3 - Results and discussion

3.1 - Optical anisotropy

The polarized Raman spectra of $BiSr_2CaCu_2O_7$ present strong anisotropy along the different axes of the crystal. Strong bands are observed at 627 and 91 cm⁻¹ along the ZZ geometry and they are nearly absent in the XX and YY configurations (Fig.1). On the contrary, in the latter two geometries, strong features are apparent at 582 and 554 cm⁻¹.

As we already mentionned in the case of Bi_{2212} and Bi_{2223} and their lead doped compounds this anisotropy can help in the determination of the orientation of any surface or thin film^(9,10).

3.2 - Tetragonal structure of Bi₁₂₁₂

If the ZZ polarized spectrum of Bi₁₂₁₂ is quite different from the XX and YY spectra, the latter two are nearly identical. This result is similar to that we already observed for $T\ell Ba_2CaCu_2O_7$ ($T\ell_{1212}$)⁽¹¹⁻¹³⁾. This is in agreement with a tetragonal structure P4/mmm for both materials. In such a case, when passing from Bi₂₂₁₂ into Bi₁₂₁₂, there is a reduction of the two formula unit cell into one formula unit cell with a $\frac{1}{2}(a+b)$ gliding. In particular, the pyramids which are at staggered positions in Bi₂₂₁₂ become in mirror image in Bi₁₂₁₂ (Fig.2). Note that the figure 2 represents only ideal limiting structures.

3.3 - Bond ordering in Bi₁₂₁₂

It is well known that in the Raman spectra of Bi₂₂₁₂ as well as of Bi₂₂₂₃, the (Bi)O(Cu) stretching is represented by a doublet around 650 and 625 cm⁻¹. This doublet is due to the in-phase and out-of-phase O(Bi) unharmonic motions in two close, staggered and antiparallel O-Bi bonds along the c-axis ; these oxygen atoms occupying the apical position of the pyramids. The band at 650 cm⁻¹ appears in all structures where the pyramidal networds are in staggered position⁽¹⁴⁻¹⁷⁾ and where the antiparallel bonds are close and identical. We can call it "staggered (Bi)O(Cu) motion". Note that the structural modulation incommensurability, which is well evidenced in



Fig.2 - Scheme of the spatial structure of Bi₁₂₁₂ and Bi₂₂₁₂.

TABLE 1 - PHONONS WAVENUMBERS OF BiSr₂CaCuO₂O₇ AND TlBa₂CaCu₂O₇.

	BiSr ₂ CaCu ₂ O ₇				TℓBa ₂ CaCu ₂ O ₇			
	Wave- number cm ⁻¹	Scatterin geometry	Symmetry & Assignmen	Wa nur cr	ave- nber n ⁻¹	Scatterin geometry	Symmetry & Assignment	
	627	ZZ	A _{1g} (Bi)O(Cu)		518	ZZ	A _{1g} (Tl)O(Cu)	
	519 333 273* 170 132 91	ZZ ZZ ZZ	A_{1g} (Cu)O _z in phase		395 80*	ZZ ZZ	A_{1g} (Cu) O_z in phase	
		ZZ ZZ ZZ	A _{1g} Cu _z		143	ZZ	A_{1g} Cu_z	
	67	ZZ	A _{1g} Sr _z		116	ZZ	A _{1g} Ba _z	
	582 554 469 305 263	XX,YY XX,YY XX,YY XX,YY XX,YY	$\begin{array}{c} O(Sr)\\ O(Bi,Sr)\\ B_{1g} (Cu)O_z \text{ out of}\\ phase \end{array}$		516 400 282 44*	XX,YY XX,YY XX,YY XX,YY XX,YY	$\begin{array}{c} O(Ba) \\ O(T\ell,Ba) \\ B_{1g} (Cu)O_z \text{ out of} \\ phase \end{array}$	
	144 50	XX,YY XX,YY			.16*	ХХ,ҮҮ		

*component from the other geometry.

these materials as having strong consequences on the properties out of the Brillouin zone center, has no direct effect on our Raman spectrum.

Such a staggered motion disappears in the Raman spectrum of Bi_{1212} where only one (Bi)O(Cu) band is observed at 627 cm⁻¹ which represents the motion of oxygen beween the Bi and Cu atoms on the same *c*-axis.

We want to mention that the same bondordering with only one single (Bi)O(Cu) band is also observed in $I.Bi_{2212}$ where the structure is similarly reduced to only one formula unit cell⁽¹⁸⁾ and where the O-Bi bonds become far apart from one to another.

In the case of lead-doped Bi_{2212} , the mechanical coupling also disappears as the antiparallel bonds are no longer identical and therefore only one O(Bi) stretching band is observed along the ZZ scattering symmetry.

3.4 - Other spectra-structure correlations

In one formula unit cell tetragonal structure with space group P4/mmm, the irreductible representation of Bi_{1212} is⁽¹⁹⁻²¹⁾:

 $\mathbf{4A_{1g}}+\mathbf{B_{1g}}+\mathbf{5E_g}+\mathbf{7A_{2u}}+\mathbf{B_{2u}}+\mathbf{8E_u}$

where :

$4A_{1g} + B_{1g} + 5E_{g}$	are Raman active are infrared active			
$6A_{2u} + 7E_{u}$				
B_{2u}	is silent.			

The A_{1g} modes include the motions of Ba, Cu, O_1 and O_2 atoms while the B_{1g} modes represent O_2

motions and the E_g modes involve the motions of Ba, Cu, O₁, O₂ atoms.

The ZZ polarized spectrum gives all four vibrations expected for A_{1g} symmetry, which represent the motions of Oxygen, Cu and Sr along the c-axis respectively at 627, 132 and 67 cm⁻¹ and in-phase oxygen motion in the Cu-O plane at 333 cm⁻¹. Note that the motion of Bi is not active.

From the XX and YY geometries, by comparison with the Raman spectrum of $T\ell_{1212}$, the oxygen (Sr) motion along the *c*-axis can be located at 554 cm⁻¹ and the out-of-phase oxygen motion in the Cu-O plane, with B_{1g} symmetry, is found at 305 cm⁻¹. All remaining phonons can be assigned to components of E_g symmetry although the E_g modes may totally appear in the ZX and ZY scattering geometries.

The correspondance of the phonons of

BiSi₂CaCu₂O₇ and their a ssignment with those of $T\ell Ba_2CaCu_2O_7$ are displayed in Table 1. As also indicated in the Table, some unassigned weak features could be due to a component from the other scattering symmetry.

4 - Conclusion

BiSr₂CaCu₂O₇ superconductor is prepared and its polarized Raman recorded and assigned by comparison with those of isostructural $T\ell Ba_2CaCu_2O_7$. High optical anisotropy is found between the c-axis and the other axes. The identity of XX and YY polarized spectra is in agreement with a tetragonal structure. The (Bi)O(Cu) bond homogeneity along the c-axis in Bi₁₂₁₂ is found and contrasts with staggered and antiparallel (Bi)O bonds in Bi₂₂₁₂ where in-phase and out-of-phase (Bi)O motions along the c-axis was recorded. Note that the latter coupling also disappears in doped (Bi/Pb)₂₂₁₂ where the staggered bonds Bi-O and Pb-O are no longer identical.

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